

Analytical functions for the calculation of hyperspherical potential curves of atomic systems

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The Hyperspherical Adiabatic Approach (HAA) [1, 2, 3] is an adequate method to treat N -body systems interacting with the long-range Coulombian forces due to its molecular-like description using potential curves that bring us to mind the spirit of the Born-Oppenheimer approximation. For two electron atomic systems ($N=3$) the HS coordinates are, $R^2 = r_1^2 + r_2^2$, $\tan \alpha = r_1/r_2$. With those coordinates the hyperspherical Schrödinger equation has the compact form

$$\left[\frac{d^2}{dR^2} - \frac{\hat{U}(R; \Omega) - 1/4}{R^2} + \varepsilon \right] (\sin \alpha \cos \alpha R^{5/2} \psi(R, \Omega)) = 0, \quad (1)$$

where ε is the system energy and the operator $\hat{U}(R; \Omega)$ depends on all compact variables $\Omega = (\alpha, \phi_i, \theta_i; i = 1, 2)$ and on the hyper-radius R through the expression

$$\hat{U} = C_2(O(3N - 3)) + R\hat{V}(R; \Omega), \quad (2)$$

where C_2 is the Casimir operator of the $O(3N - 3)$ group and \hat{V} is the interparticle potential energy. The eigenstates of \hat{U} are the channel functions, $\Phi_\lambda(R; \Omega)$ with eigenvalues forming potential curves $U_\lambda(R)$ as function of R . In the Hyperspherical Adiabatic Approach (HAA) the total wave function is expanded in the basis formed by the channel functions,

$$\psi(R, \Omega) = \sum_\lambda F_\lambda(R) \Phi_\lambda(R; \Omega), \quad (3)$$

which leads to a set of coupled equations for the radial components $F_\lambda(R)$.

By introducing the variable $z = \tan(\alpha/2)$ [2, 4] we present a fast convergent analytical angular basis functions to expand $\Phi_\lambda(R; \Omega)$ which generalizes the Jacobi polynomials. We show that these functions, obtained by selecting the diagonal terms of the angular equation with respect to the total angular momentum function $\mathcal{Y}_{\ell_1 \ell_2}^{LM}(\hat{r}_1, \hat{r}_2)$, allow efficient diagonalization of the Hamiltonian for all values of the hyperspherical radius [5].

For the Li^+ ion and helium atom with $L = 0$ and total spin $S = 0$ the ground state energy obtained as a function of the number of radial coupled components are given on table I.

Table 1: Ground state energy of the Li^+ as a function of the number of coupled radial equations.

N_c	Li^+ Energy (a.u.)	Error (ppm)	He Energy (a.u.)	Error
1	-7.262 640 06	2372.737	-2.895 553 56	2 813.911
3	-7.279 704 90	28.639	-2.903 611 47	38.882
7	-7.279 897 74	2.149	-2.903 717 07	2.515
13	-7.279 909 25	0.568	-2.903 723 00	0.475
21	-7.279 910 83	0.351		

The method also allows the calculation of excited states when the correct radial asymptotic condition is considered [6]. Some results are shown on Table II.

Table 2: Lowest binding energies of the Li^+ [5] and Helium [7]

State	Li^+ Energy (a.u.)	Error (ppm)	He Energy (a.u.)	Error (ppm)
1s1s	-7.279 910 8	0.35	-2.903 723 00	0.48
1s2s	-5.040 865 9	2.15	-2.145 956 96	7.96
1s3s	-4.733 725 0	6.44	-2.061 242 87	14.13
1s4s	-4.629 749 1	7.15	-2.033 567 78	9.31
1s5s	-4.582 401 5	5.64	-2.021 164 76	5.98

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